

目 录

1 基本的状态方程	2
1.1EOS_JWL	2
2.2EOS GRUNEISEN	2
2.3EOS_LINEAR_POLYNOMIAL	3
2.材料模型.....	3
2.1MAT_HIGH_EXPLOSIVE_BURN	3
RDX.....	5
HMX	5
TNT	5
1.2MAT_NULL	6
空气	6
水.....	7
1.3MAT JOHNSON COOK	7
紫铜.....	8
钢	8
1.4 MAT_PLASTIC_KINEMATIC	9
钢.....	10
高导无氧铜.....	10
土壤	10
1.5MAT_STEINBERG	10
高导无氧铜.....	12
1.6MAT_JOHNSON_HOLMQUIST_CERAMICS	12
B4C陶瓷	14
1.7MAT JOHNSON HOLMQUIST CONCRETE	14
混凝土.....	14
3 其它材料参数.....	15
LY12CZ 铝合金	15

主要材料模型及参数

1 基本的状态方程

1.1 EOS_JWL

	1	3	5	#				
Vathte	EOD	A	H	风1	M1	FF	E	vo
Ty=	1	F	干	F	P	F	F	F

Remarks

The JWL equation of state defines the pressure as

$$P = A \left(1 - \frac{\omega}{R, V} \right) e^{-\beta, V} + B \left(1 - \frac{\omega}{R, V} \right) e^{-\beta, V} + \frac{\omega E}{V}$$

and is usually used for detonation products of high explosives.

2.2 EOS_GRUNEISEN

	2	4	4	s	省	T		
VattNe	to[D	c	Si	2	\$3	GAHAO	A	Ea
Iyie			F	P	F	F	F	P

Caidl

Vatb	Vo							
Tyr	F							

Remarkst

The Gruoeisen equation of state with cubic shock velocity-particle velocity defines pressure for compressed materials as

$$P = \frac{\rho_0 C^2 \mu \left[1 + (1 - \frac{1}{\mu}) \mu - \frac{1}{\mu} \right]}{\left[1 - (S_1 - 1) \mu - S_2 \frac{\mu^2}{\mu - 1} - S_3 \frac{\mu^3}{(\mu - 1)^2} \right]^2} + (\gamma_0 + \mu \mu) E$$

ant fur expanded materials as

$$P = AC' \mu + (X\% + \mu)E$$

where C is the intercept of the γ_p curve: S_x , S_z , and S_y are the coefficients of the slope of the γ_p curve: T_0 is the Grüneisen parameter μ is the first order volume correction to ρ_0 and $\mu = \frac{\rho}{\rho_0} - 1$.

2.3 EOS_LINEAR_POLYNOMIAL

(对 EOS_GRUNEISEN 进行线性化)

Can1	1	2	J	1	5	6		
Veiahls	EOSD	CO	C1	C	C3	C4	C5	Cn
Trp	1		F	扩	节	F	F	

Can2								
Vataht	BP	w						
Typ		F						

Remarks

1. The linear polynomial equation of state is linear in internal energy. The pressure is given by;

$$P=C+C\mu+G\mu^2+G\mu^1+(G+C\mu+C,\mu)E,$$

where terms $C2u?$ and Con^2 are set to zero if $\mu < 0$, $\mu = \frac{\rho}{\rho_0} - 1$ and $\frac{\rho}{\rho_0}$ is the ratio of current density to initial density.

2. 材料模型

2.1 MAT_HIGH_EXPLOSIVE_BURN

Can1	Parnat							
Can1	1	2	J	4	5		1	
Vwiahia	MID	RO	D	PCJ	BETA	K	G	SIOY
Type	1	F	F	F	F	卡	F	F

Burn fractions, F, which multiply the equations of states for high explosives, control the release of chemical energy for simulating detonations. At any time, the pressure in a high explosive element is given by:

$$p = Fp(V, E)$$

where p is the pressure from the equation of state (either types 2 or 3), V is the relative volume, and E is the internal energy density per unit initial volume.

In the initialization phase, a lighting time t is computed for each element by dividing the distance from the detonation point to the center of the element by the detonation velocity D . If multiple detonation points are defined, the closest detonation point determines t . The burn fraction F is taken as the maximum

where

$$F_1 = \begin{cases} \frac{2(t-t_i)DA_{um}}{3v_r} & \text{if } t > t_i \\ 0 & \text{if } t \leq t_i \end{cases}$$

$$F_2 = \beta = \frac{1-V}{1-V_{cj}}$$

where V is the Chapman-Jouguet relative volume and t is current time. If F exceeds 1, it is reset to 1. This calculation of the burn fraction usually requires several time steps for F to reach unity, thereby spreading the burn front over several elements. After reaching unity, F is held constant. This burn fraction calculation is based on work by Wilkins [1964] and is also discussed by Giroux [1973].

If the beta burn option is used, $BETA=1.0$, any volumetric compression will cause detonation and

$$F = F_2$$

and $\dot{\epsilon}$ is not computed.

If programmed burn is used, $BETA=2.0$, the explosive model will behave as an elastic perfectly plastic material if the bulk modulus, shear modulus, and yield stress are defined. Therefore, with this option the explosive material can compress without causing detonation.

As an option, the high explosive material can behave as an elastic perfectly-plastic solid prior to detonation. In this case we update the stress tensor to an elastic trial stress, ${}^*s_{ij}^{n+1}$:

$${}^*s_{ij}^{n+1} = s_{ij}^n + s_{ip} \Omega_{pj} + s_{jp} \Omega_{pi} + 2G \dot{\epsilon}_{ij}^n dt$$

where G is the shear modulus, and $\dot{\epsilon}$ is the deviatoric strain rate. The von Mises yield condition is given by:

$$\phi = J_2 - \frac{\sigma_y^2}{3}$$

where the second stress invariant, J_2 , is defined in terms of the deviatoric stress components as

$$J_2 = \frac{1}{2} s_{ij} s_{ij}$$

and the yield stress is σ_y . If yielding has occurred, i.e., $\phi > 0$, the deviatoric trial stress is scaled to obtain the final deviatoric stress at time $n+1$:

$$s_{ij}^{n+1} = \frac{\sigma_y}{\sqrt{3J_2}} {}^*s_{ij}^{n+1}$$

If $\phi \leq 0$, then

$$s_{ij}^{n+1} = {}^*s_{ij}^{n+1}$$

Before detonation pressure is given by the expression

$$p^{n+1} = K \left(\frac{1}{V^{n+1}} - 1 \right)$$

where K is the bulk modulus. Once the explosive material detonates:

$$x_{ij}^{n+1} = 0$$

and the material behaves like a gas.

RDX

密度: 1.69E+3 kg/m³: D:8310m/s: Pci:30.45 G

A:850 Gpa: B:18 Gpa: R₁:4.6: R₂:13: w₀:38: E₀:10MJ/kg

For (g-cm-us) :

*MAT_HIGH_EXPLOSIVE_BURN

1 1.69 8.310 0.3015 0

*EOS_JWL

8.50 0.18 4.6 1.3 0.38 10 e-02 1.00

HMX

密度: 1.891E+3 kg/m³ D:9910m/s, Pcj:42Gpa

: aBp:7.31 Gpa; R₁:4.1; R₂:1.00; w₀:30; E₀:10.5 M

For (g-cm-us) :

*MAT_HIGH_EXPLOSIVE_BURN

1 1.89 9.910 0.42 0

*EOS_JWL

1 7.783 0.071

4.2 1.0 0.30 10.5 e-02 1.00

TNT

密度: 1.63 E+3 kg/m³ D:6930 m/s; Pci:27 Gpa:

A:371.2 Gpa; B:3.21 Gpa; R:4.1 5; R₂:0.9 5; w₀:30 E₀:29 MJ/kg

For (g-cm-us) :

*MAT	HIGH	EXPLOSIVE	BURN						
	1	1.63	6.930	0.27	0				
*EOS	JWL								
	1	3.713	0.0743	4.15	0.95	0.30	7.0 e-02	1.00	

1.2MAT_NULL

Cmd1	2	3	3	6				
Vaniahe	MID	RO	Pe	MU	1EROO	CFROD	VM	PR
Tba*	1	F	F	F	P	F	F	
Dfangs		p	0. a	40	0北	.0	核林	y.

The null material must be used with an equation of-state. Pressure cutoff is negative in tension.
A(deviatoric)viscous stress of the form

$$\sigma = \mu e'$$

$$\left[\frac{N}{m^2} \right] \approx \left[\frac{N}{m^2 \cdot s} \right] \left[\frac{1}{s} \right]$$

is computed for nonzero μ where e' is the deviatoric strain rate. μ is the dynamic viscosity with unit of [Pascal*second].

- The null material has no shear stiffness and hourglass control must be used with great care. In some applications, the default hourglass coefficient might lead to significant energy losses. In general for fluid(s), the hourglass coefficient QM should be small (in the range 1.0E-4 to 1.0E-6 in the SI unit system for the standard default IHQ choice).
- The Null material has no yield strength and behaves in a fluid-like manner.
- The pressure cut-off, PC, must be defined to allow for a material to "numerically" cavitate. In other words, when a material undergoes dilatation above certain magnitude, it should no longer be able to resist this dilatation. Since dilatation stress or pressure is negative, setting PC limit to a very small negative number would allow for the material to cavitate once the pressure in the material goes below this negative value.

空气

***MAT NULL**

RO=1.25 kg/m³ PC=-1.0pa(<0) MU=1.7450E-3 (动力粘性系数)

***EOS_LINEAR_POLYNOMIAL**

1,0 Gpa, 0Gpa, 0 Gpa, 0, 0.4, 0.4

253312.5, 1.0

For (g-cm-us) :

*MAT_NULL							
	30.125e-02						
		-1.0E-12	1.749E-7	00000	0000	00	00
*EOS GRUNEISEN							
	3	0.3444					
			00000	0000	00000	1.40	00
							00

```

00
/* EOS LINEAR POLYNOMIAL
3 0
0 0 0 0.4 0.4 0
2.5000E-6 1

```

水

***MAT_NULL**

1. R0=998.21kg/m³, PC=-10.0 pa, MU=0.8684E-3 0, 0, 0, 0

C:1480m/s, SL:2.5 83 0.2268 0.4934 A:0.27 E0:Q

Y0:1

For (g-cm-us) :

```

*MAT_NULL
1 0.998
-1.0E-11 0.8684E-5 00000 0000 00 00
*EOS GRUNEISEN
1 1.65
00 1.92 -0.096 00000 0.350 00 00

```

1.3MAT_JOHNSON_COOK

Cant.1 1 1 1 4 S T

Vsise	MtD	Rn	a	五	FR	DTF	VP	
Dm	l	开	F	F	F	F	护	
Defadt	mmc	*	si		nime	0.0	0.0	

Cant 2

WaiHe	A	b	N	c	M	TM	TR	EYSO
Tye	F	F	F	F	F	F	F	F
Delt	me	0,0	0.0	0.0	nooe	none	mne	be

Cand3

VmtsHe	CF	PC	SFAL. 1.	π	Di	2	33	D4
--------	----	----	----------	---	----	---	----	----

Pye	F	g	F	F	F	F	g	F
Dasdt	ncne	0,0	2.0	0.0	0.0	0.0	0.0	

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